

EPR of Gd^{3+} in Hydrated Double Nitrate Single Crystals

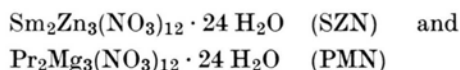
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EPR of Gd^{3+} in $Sm_2Zn_3(NO_3)_{12} \cdot 24 H_2O$ and $Pr_2Mg_3(NO_3)_{12} \cdot 24 H_2O$ single crystals has been studied at X-band. In addition to the allowed fine structure lines ($\Delta M = \pm 1$) some weak low field lines identified as $\Delta M = \pm 2$ transitions, have been observed.

The electron paramagnetic resonance (EPR) of Gd^{3+} has been previously studied in single crystals of hydrated double nitrates, $M_2'''M_3''(NO_3)_{12} \cdot 24 H_2O$ ($M''' = La, Bi$ and $M'' = Mg, Zn$) [1, 2] and in $Pr_2Zn_3(NO_3)_{12} \cdot 24 H_2O$ [3]. However, the weak lines which occur at the low magnetic field side of the allowed ($\Delta M = \pm 1$) fine structure lines in the EPR spectra of Gd^{3+} in double nitrates have not been investigated previously. As will be shown in this paper these lines have been identified as arising from the higher order EPR transitions ($|\Delta M| > 1$). This paper describes the EPR of Gd^{3+} in



single crystals from 300 K to 77 K.

The crystal structure of $Ce_2Mg_3(NO_3)_{12} \cdot 24 H_2O$ has been determined by Zalkin et al. [4]. The SZN and the PMN can be expected to have similar structure. The primitive cell containing one formula unit is rhombohedral with space group $R\bar{3}$ [4, 5]. The unit cell contains three divalent ions situated at two different lattice sites with point symmetry C_3 and C_1 . The trivalent ions are found at sites of C_1 point symmetry and the rest of the atoms are in positions of the general type. Culvahouse et al. [6] have shown that there are two trivalent ions along the long diagonal of the rhombohedron and 5.48 \AA to either side of the centre of the cell. The long diagonal of the rhombohedron is along the trigonal axis. Each trivalent ion is coordinated with 12 oxygens belonging to six nitrate ions, located at the corners of a somewhat irregular icosahedron.

Single crystals of SZN and PMN doped with Gd^{3+} were grown at room temperature by slow

evaporation of aqueous solutions of $M'''(NO_3)_3 \cdot 6 H_2O$ ($M''' = Sm, Pr$) and $M''(NO_3)_2 \cdot 6 H_2O$ ($M'' = Zn, Mg$) mixed in stoichiometric ratios. The Gd^{3+} was introduced into the host lattices by adding a small amount (0.2% by weight) of $Gd(NO_3)_3 \cdot 6 H_2O$. The crystals grow in flat hexagonal plates, the plane of which is perpendicular to the trigonal axis [7]. A Varian V-4502 X-band EPR spectrometer with a 9-inch magnet and 100 kHz field modulation was used to record the spectra. As a reference for the magnetic field strength, the resonance line of DPPH with $g = 2.0036$ was used.

For an arbitrary orientation of the crystal, the EPR spectrum consists of a single set of seven lines between $\sim 0.25 \text{ T}$ and $\sim 0.43 \text{ T}$ which arise from the allowed fine structure transitions ($\Delta M = \pm 1$) of the Gd^{3+} centre (formed by the substitution of M''' by Gd^{3+}). In addition to these lines there are many weak lines below $\sim 0.24 \text{ T}$ (Figure 1). Angular variation studies of the spectrum reveal the following:

1. The maximum spread of the spectrum (z axis) corresponding to $\Delta M = \pm 1$ transitions occurs when the external magnetic field is parallel to the trigonal axis (c axis) of the crystal. This axis can be located accurately from the external morphology of the crystal.
2. The subsidiary maxima (x axis) occur when the external magnetic field is approximately perpendicular to the c axis. However, when the field is at 90° to c axis, all the observed transitions do not attain an extremum. The departure from

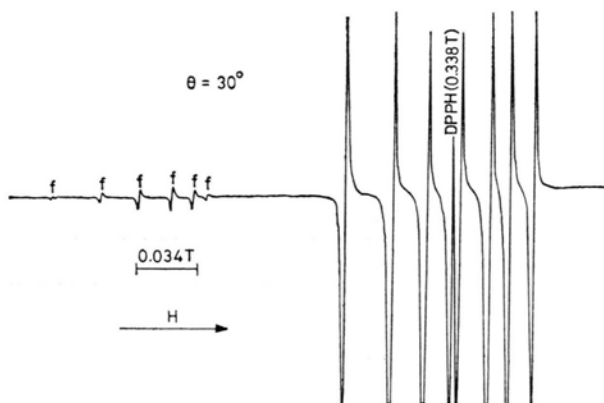


Fig. 1. EPR spectrum of Gd^{3+} in $Sm_2Zn_3(NO_3)_{12} \cdot 24 H_2O$ single crystals at 300 K; with H makes an angle 30° to the z axis in the z plane. Forbidden fine structure lines are marked f.

the perpendicular direction is as much as $\sim 6^\circ$ and the corresponding change in field position is ~ 0.001 T which is more than the experimental error.

3. The spectrum shows a periodicity of $\pi/3$, when the crystal is rotated in the plane perpendicular to the c axis. However, in this plane the angular variation of line positions is very small, therefore the site symmetry can be treated as being practically trigonal.

The spin-Hamiltonian for Gd^{3+} corresponding to a trigonal symmetry and with the z axis parallel to the c axis can be written as [8]

$$\begin{aligned} \mathcal{H} = g \beta S \cdot H &+ \frac{1}{3} b_2^0 O_2^0 + \frac{1}{60} b_4^0 O_4^0 \\ &+ \frac{1}{1260} b_6^0 O_6^0 + \frac{1}{3} b_4^3 O_4^3 \\ &+ \frac{1}{36} b_6^3 O_6^3 + \frac{1}{1260} b_6^6 O_6^6 \end{aligned} \quad (1)$$

where the symbols have their usual meaning. The crystal field parameters b_4^3 and b_6^3 are identically zero for C_{3h} and D_{3d} symmetry and not zero for C_{3v} symmetry [2].

To find the parameters of the spin-Hamiltonian (1) we have used the approximate equations for $H \parallel z$ given in [1]. With these equations we determine the parameters b_2^0 , b_4^0 , b_6^0 and g_{\parallel} . The absolute value of the constant b_6^6 was found from the variation of the position of lines at $H \perp z$ axis. The parameters for Gd^{3+} in SZN and PMN at 300 K and 77 K are given in Table 1. The signs of the parameters b_n^m given in this table are only relative and have been determined assuming b_2^0 to be positive.

Table 1. Spin-Hamiltonian parameters for Gd^{3+} in $\text{Sm}_2\text{Zn}_3(\text{NO}_3)_{12} \cdot 24 \text{H}_2\text{O}$ and $\text{Pr}_2\text{Mg}_3(\text{NO}_3)_{12} \cdot 24 \text{H}_2\text{O}$ at 300 K and 77 K. The b_n^m parameters are given in units of 10^{-4} cm^{-1} .

Spin-Hamiltonian parameters	$\text{Sm}_2\text{Zn}_3(\text{NO}_3)_{12} \cdot 24 \text{H}_2\text{O}$		$\text{Pr}_2\text{Mg}_3(\text{NO}_3)_{12} \cdot 24 \text{H}_2\text{O}$	
	300 K	77 K	300 K	77 K
b_2^0	126.6 ± 0.5	137.6 ± 0.5	113.7 ± 0.5	123.2 ± 0.5
b_4^0	0.58 ± 0.05	0.65 ± 0.05	0.29 ± 0.05	0.48 ± 0.05
b_6^0	0.48 ± 0.05	0.74 ± 0.05	0.44 ± 0.05	0.79 ± 0.05
b_6^6	5.6 ± 1	5.6 ± 1	4.6 ± 1	5.6 ± 1
g_{\parallel}	1.9917 ± 0.0005	1.9917 ± 0.0005	1.9922 ± 0.0005	1.9922 ± 0.0005
g_{\perp}	1.992 ± 0.002	1.992 ± 0.002	1.995 ± 0.003	1.995 ± 0.003

To identify the observed low field lines we have investigated higher order EPR transitions ($|\Delta M| > 1$) for the Gd^{3+} centre. For this purpose we have evaluated the positions of all the fine structure transitions. The approximate theoretical expressions for the angular variation of line positions were obtained by considering the predominant $b_2^0 O_2^0$

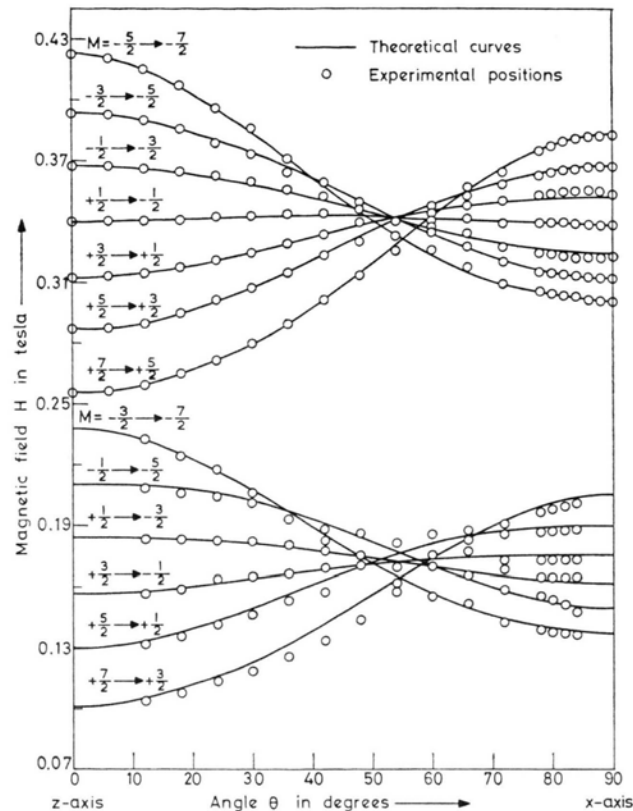


Fig. 2. Angular dependence of Gd^{3+} EPR lines in the zx plane for $\text{Sm}_2\text{Zn}_3(\text{NO}_3)_{12} \cdot 24 \text{H}_2\text{O}$ at 300 K.

term up to second order of perturbation and the rest of the terms of the spin-Hamiltonian only up to first order. Figure 2 shows the angular variation of the spectrum in the zx plane. In the figure Θ is the angle between the external magnetic field and the z axis. The reasonably good agreement of the calculated angular variation with the experimental

one suggests that the weak lines at the low magnetic field side of the allowed ($\Delta M = \pm 1$) fine structure lines (Fig. 1) are due to $\Delta M = \pm 2$ transitions.

Acknowledgements

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- [1] R. S. Trenam, Proc. Phys. Soc. London **A 66**, 118 (1953).
- [2] H. A. Buckmaster, J. C. Dering, and J. I. Fry, J. Phys. **C 1**, 599 (1968).
- [3] V. K. Jain and T. M. Srinivasan, Pramana **10**, 155 (1978).
- [4] A. Zalkin, J. D. Forrester, and D. H. Templeton, J. Chem. Phys. **39**, 2881 (1963).
- [5] M. R. Anderson, G. J. Jenkin, and J. W. White, Acta Cryst. **B 33**, 3933 (1977).
- [6] J. W. Culvahouse, W. Unruh, and R. C. Sapp, Phys. Rev. **121**, 1370 (1961).
- [7] R. A. Fisher, Rev. Sci. Instrum. **43**, 386 (1972).
- [8] H. A. Buckmaster and Y. H. Shing, Phys. Stat. Sol. (a) **12**, 325 (1972).